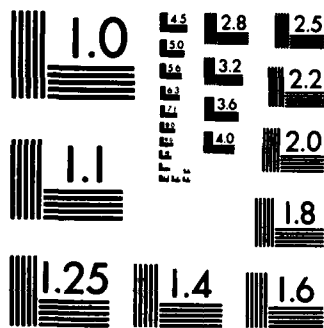


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006 00 DAVID COLTON

007 INVERSE SCATTERING PROBLEM FOR TIME-HARMONIC WAVES 00

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015 THE INVERSE SCATTERING PROBLEM FOR TIME-HARMONIC  
016 ACOUSTIC WAVES\*

017 DAVID COLTON†

018 **Abstract.** The inverse scattering problem we are considering in this paper is to determine the shape of a  
019 sound-soft, bounded, connected obstacle from a knowledge of the time-harmonic incident wave with frequency  
020 in the resonant region and the far field pattern of the scattered wave. After some introductory remarks, we begin  
021 by describing the method of integral equations and the null-field method for solving the direct scattering  
022 problem. This enables us to define an operator  $T$  mapping the boundary of the scattering obstacle and the  
023 incident field onto the far field pattern. The inverse scattering problem is to invert this operator. In order to do  
024 this, we first must examine the range of  $T$ , i.e. to characterize the class of far field patterns, and to establish the  
025 existence of  $T^{-1}$  on the range of  $T$ , i.e. to show the uniqueness of the solution to the inverse scattering problem.  
026 This analysis shows that for a given measured far field pattern, in general no solution exists to the inverse  
027 scattering problem, and if a solution does exist, it does not depend continuously on the measured data, i.e. the  
028 problem is improperly posed. This motivates us to consider a linearized model and to examine various methods  
029 for studying linearized improperly posed problems based on the ideas of a priori assumptions and compactness  
030 arguments. We then consider a simple model problem that focuses on the nonlinear character of the inverse  
031 scattering problem and, motivated by our study of the linearized model, reformulate the inverse scattering  
032 problem as a problem in constrained optimization. We conclude by considering the numerical solution of this  
033 nonlinear optimization problem.

034 "Approach your problems from the right end and begin with the answers. Then, one day,  
035 perhaps you will find the final question." from "The Hermit Clad in Crane Feathers" in *The Chinese*  
036 *Maze Murders*, by R. Van Gulik.

037 **1. Introduction.** This paper is devoted to surveying some of the recent developments  
038 associated with the inverse scattering problem for acoustic waves. This problem is perhaps  
039 the most famous of the inverse and improperly posed problems arising in mathematical  
040 physics; indeed, the name itself is "improperly posed" in the sense that there are  
041 numerous "inverse scattering problems" arising in the theory of acoustic wave propaga-  
042 tion. Hence it is necessary to clarify at the beginning exactly which of these is to be  
043 considered. To get a glimpse of the various types of inverse scattering problems that can  
044 arise in applications, it is instructive to first briefly consider some typical problems  
045 associated with acoustic scattering. We shall later be more precise concerning the  
046 physical derivation and mathematical modeling of both the direct and inverse problem,  
047 and hence, for the moment, we shall content ourselves with some simple heuristic  
048 descriptions. Consider an acoustic wave propagating in a homogeneous, isotropic  
049 medium. In the absence of any inhomogeneities, the wave will continue to propagate and  
050 nothing of physical interest will happen. However, if there are inhomogeneities present,  
051 then the wave will be "scattered" or "diffracted" and we can express the total field as the  
052 sum of the original "incident" wave and the "scattered" wave. The behavior of the  
053 scattered wave will depend on both the incident wave and the nature of the inhomogene-  
054 ities in the medium, and this in turn is reflected in the mathematical model, e.g. what is  
055 the incident wave, is the inhomogeneous region connected, are the boundaries of the  
056 inhomogeneous medium bounded or infinite, what type of boundary conditions are  
057 appropriate, etc.? The direct problem is, given this information, to find the scattered wave  
058 and in particular, its behavior at large distances from the inhomogeneities, i.e. its "far  
059 field" behavior. The inverse problem takes this answer to the direct scattering problem as  
060 its starting point and asks what is the nature of the inhomogeneities which gave rise to  
061 such a far field behavior.

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Dense Sets and Far Field Patterns  
in Electromagnetic Wave Propagation

by

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# Abstract

It is shown that the electric far field patterns corresponding to the scattering of entire incident fields by a bounded perfectly conducting obstacle are dense in the space of square integrable tangential vector fields defined on the boundary of the unit sphere if and only if there does not exist a Maxwell eigenfunction that is an electromagnetic Herglotz pair, i.e. a solution  $\{E, H\}$  of Maxwell's equations defined in all of space such that

$$\lim_{r \rightarrow \infty} \frac{1}{r} \iint_{|x| < r} (|E(x)|^2 + |H(x)|^2) dx < \infty .$$



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and Elastic Waves*

Far Field Patterns in Acoustic and  
Electromagnetic Scattering Theory

Abstract

A basic task in the investigation of the inverse scattering problem for time-harmonic acoustic and electromagnetic waves is the study of the class of far field patterns corresponding to the scattering of entire incident fields of a given wave number by a bounded obstacle. Indeed if  $\mathcal{T}$  denotes the operator mapping the incident field and scattering obstacle onto the far field pattern, then the inverse scattering problem is to construct  $\mathcal{T}^{-1}$  defined on the range of  $\mathcal{T}$ , and the determination of this range is nothing more than the description of the class of far field patterns. Unfortunately, little is known concerning this class except for the well known fact that the far field patterns are entire functions of their independent (complex) variables for each positive fixed value of the wave number ( [3] ), i.e. the range of  $\mathcal{T}$  is not all of  $L^2(\Omega)$  where  $\Omega$  is the unit sphere. We note that this implies that the inverse scattering problem is an improperly posed problem since the far field patterns are in practice determined from inexact measurements.

Recently Colton ( [1] ) and Colton and Kirsch ( [2] ) have investigated the case of acoustic scattering and asked the question if the class of far field patterns corresponding to a fixed scattering obstacle and all entire incident fields is dense in  $L^2(\Omega)$ . The rather surprising answer to this question is that if the impedance of the scattering obstacle is positive, then the far field patterns are dense in  $L^2(\Omega)$ , whereas if the scattering obstacle is sound-soft or sound-hard then the far field patterns are dense in  $L^2(\Omega)$  if

## Uniqueness Theorems for the Inverse Problem of Acoustic Scattering

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Uniqueness theorems are obtained for the problem of determining the shape of a sound-soft or sound-hard obstacle from a knowledge of (1) the far-field pattern at a fixed value of the wave number and a finite number of distinct incident fields, or (2) the total scattering cross section for an interval of wave numbers and the incident field propagating in an arbitrary direction.

### 1. Introduction

THE INVERSE SCATTERING PROBLEM for acoustic waves forms the basis of a wide variety of areas in the engineering sciences involving remote sensing and imaging, and for this reason has been the object of intensive study by scientists in a number of diverse disciplines. Since around 1970 progress has been particularly rapid, and for a survey of these recent results we refer the reader to the expository papers by Colton (1983) and Sleeman (1982). However, in this intensive and prolonged effort there are at present only a small number of results available on the uniqueness of the solution to the inverse scattering problem. The purpose of this paper is to add several additional uniqueness theorems to this sparse collection, our motivation coming from some recent numerical results of Andreas Kirsch who considers inverse scattering problems not covered by previously known uniqueness theorems and which in fact seem to exhibit non-uniqueness (Kirsch, 1982). However, before we can describe our results we must be more precise as to what inverse scattering problem we are considering, since the term "inverse scattering problem" is not uniquely defined. To this end we consider a plane time-harmonic acoustic wave moving in the direction  $\alpha$  that is scattered by a bounded connected domain  $D$  in  $\mathbb{R}^3$  which is assumed to be either "sound-soft" or "sound-hard". Then if we factor out the periodic dependence on time and denote the total field by  $u(x)$ ,  $x \in \mathbb{R}^3 \setminus D$ , the scattered field by  $u^s(x)$ , and the (positive) wave number by  $k$ , we have that

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Uniqueness of Solutions to the Inverse  
Acoustic Scattering Problem

David Colton and B D Sleeman

1. INTRODUCTION

The inverse acoustic scattering problem forms the basis of a wide range of problems in the engineering sciences including for example remote sensing and imaging and consequently has been the object of intensive study in recent years. For an overview of recent contributions we cite the expository articles [1] and [6]. Despite this intensive research there are only a few results concerned with the question of uniqueness of solutions to the inverse scattering problem. In this paper we review known uniqueness results and report on some new developments. Full proofs of the new results are to be found in [3].

Before we can adequately describe what is meant by the inverse scattering problem it is necessary to recall some fundamental notions concerning the direct problem. To this end we consider a plane time-harmonic acoustic wave moving in the direction  $\underline{\alpha}$  that is scattered by a bounded connected domain (the scattering obstacle)  $D$  in  $\mathbb{R}^3$  assumed to be either 'sound-soft' or 'sound-hard'. If we suppress the assumed harmonic time dependence and denote the total field by  $u(\underline{x})$ ,  $\underline{x} \in \mathbb{R}^3 \setminus D$ , the scattered field by  $u^s(\underline{x})$  and the (positive) wave number by  $k$  then  $u(\underline{x}) = e^{ik\underline{x} \cdot \underline{\alpha}} + u^s(\underline{x})$ ,  $|\underline{\alpha}| = 1$ , must satisfy the following boundary value problem for the Helmholtz equation

$$\Delta_3 u + k^2 u = 0 \quad \text{in } \mathbb{R}^3 \setminus \bar{D} \quad (1.1)$$

$$u = 0 \quad \text{on } \partial D, \quad (1.1a)$$

$$\text{or} \quad \frac{\partial u}{\partial \nu} = 0 \quad \text{on } \partial D, \quad (1.1b)$$

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*The Strong Maximum Principle for  
the Heat Equation\**

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Abstract: A mean value theorem is derived for the heat equation and is used to prove the strong maximum principle for solutions of the heat equation.

The strong maximum principle for harmonic functions is usually arrived at by appealing to the mean value theorem (c.f. [1], p. 53). It is also of course possible to simply appeal to the Hopf maximum principle ([2]), but using sledge hammers to kill flies is generally viewed as aesthetically displeasing. In contrast to the case of harmonic functions, the only proof of the strong maximum principle for the heat equation that is known to me is to invoke Nirenberg's strong maximum principle for parabolic equations ([2]). As in the case of harmonic functions, it seems desirable to provide a direct proof of this result without having to go through the subtle comparison arguments that are employed in the more general case. The purpose of this note is to provide a proof of the strong maximum principle for the heat equation based on a mean value theorem for solutions of the heat equation which we derive below. Such an approach provides a straightforward and simple proof of the strong maximum principle which avoids most of the detailed estimates of the proof of the maximum principle for more general parabolic equations.

For the sake of simplicity we shall only consider the case of the heat equation in three space dimensions.

Theorem: Let  $u(x,t)$ ,  $x \in \mathbb{R}^3$ , satisfy the heat equation

$$\Delta_x u = u_t$$

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ANALYTIC SOLUTIONS OF THE HEAT EQUATION AND SOME FORMULAS FOR LAGUERRE  
AND HERMITE POLYNOMIALS\*

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and the second author was supported in part by NSF Grant MCS-8301842.

# Abstract

We consider analytic solutions of the heat equation  $u_{xx} + u_{yy} = u_t$  defined in a cylinder and show that any such solution can be expanded in a series of polynomial solutions to the heat equation. If we define the independent complex variables  $z$  and  $\bar{z}$  by  $z = x + iy$ ,  $\bar{z} = x - iy$ , where  $x$  and  $y$  are independent complex variables, it is shown that any real-valued analytic solution of the heat equation is uniquely determined by its values on  $\bar{z} = 0$  or  $t = 0$ . Using this result, and expressing the above mentioned polynomial solutions to the heat equation in terms of Laguerre polynomials, we obtain some generating functions for Laguerre polynomials, as well as connection formulas between products of Hermite polynomials and Laguerre polynomials of argument  $r^2 = x^2 + y^2$ . These connection formulas generalize a well known result of Feldheim.

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